Existence of parabolic orbits for the restricted three-body problem

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ABSTRACT. In this paper, we show, using a variational formulation, the existence of Parabolic or homoclinic orbits at infinity of the restricted three-body problem.

$$\ddot{z}(t) + \alpha \frac{\ddot{z}}{(z(t)^2 + r^2)^{\frac{\alpha}{2} + 1}} = 0.$$

For this, we prove the existence of a minimax critical level of functionals defined on the spaces of periodic functions H^1_{2mT} , we get a sequence $(z_{\alpha,m})_{m\in\mathbb{N}}$. By the Ascoli Theorem, we prove that $(z_{\alpha,m})$ converges to a parabolic orbit.

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1. Introduction

The restricted three-body problem belongs to a large category of problems called the Hamiltonian systems. An orbit of this problem is said to be parabolic if two of the three bodies remain bounded while the third goes to infinity with vanishing velocity. The existence of infinitely many periodic and parabolic solutions of such systems has already been largely studied in the litterature ([8]-[11], [14]-[16], [18], [22]-[24], [29]-[32], [38]). In [26], Moser treated this problem using the geometry of the Bernoulli-Shift and the symbolic dynamics. In his method, one has to prove the existence of homoclinic points, which needs the verification of the transversality of stable and unstable manifolds near the hyperbolic fixed points of the systems.

Sitnikov [35] proved the existence of solutions said to be oscillatory, i.e. which correspond to motions in which two bodies remain bounded while the third limits to infinity with zero velocity. He was based on symmetry conditions. In [1], Alekseev treated the problem in a more general point of view using the theory of Poincaré [27]. Than came R.Mc Gehee [20] who studied the application of Poincaré associated to a periodic orbit of Sitnikov's exemple. He showed that the intersection of stable and unstable manifolds to a degenerated fixed point, is an homoclinic point. And nearby we can remark the behaviour of the oscillatory solutions. This result is a generalization of the theory of Slotnik [36] dealing with stable manifolds near degenerated fixed points.

Many authors have been interested in the research of periodic orbits for N-body type problems ([14], [15], [22]-[24], [29], [38] ...) and especially those of the restricted three-body problem ([3], [5], [6], [21], [33] -[35] ...). In the last few years, and precisely from the work of Rabinowitz [28], some of them focused their works on the new and powerful tool consisting on variational methods. In fact, we can find many references based on this tool, one can cite for example [2], [4], [10], [14]- [17], [20], [21], [37] ...

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However, there have been less interest in the looking for homoclinic orbits, and lesser for the research of parabolic ones. We can cite, as examples the works of A. Bahri and P.H. Rabinowitz [5] which was revisited by H. Riahi ([29] -[31]), and those of E. Séré [32], V. Coti-Zelati and I. Ekeland [17]...

In these methods, the existence of homoclinic orbits and hyperbolic fixed points is no longer studied. We avoid by the way to look for the intersection or to prove the transversality of manifolds. Thus, we transform the geometric approach to solve this problem to a purely analytic one.

This paper is organized as follows: in section 2, we set the problem in a general frame, then we introduce the particular case studied in this work. Section 3 treats the case $0 < \alpha < 1$ and contains two steps: in the first one, we prove the existence of a series of minimax critical levels of functionals defined on the spaces of periodic functions H_{2mT}^1 : we get a sequence $(z_{\alpha,m})_{m\in\mathbb{N}}$. In the second step, using the Ascoli-Arzelà Theorem, we prove that $(z_{\alpha,m})$ converges to a parabolic orbit. Finally, in section 4, taking a sequence $(\alpha_n)_{n\in\mathbb{N}} \subset]0, 1[$, and making it converge to 1, we obtain the response for $\alpha = 1$.

2. Problem setting

In this paper, we study the existence of parabolic solutions of the restricted threebody problem. We are interested in the configuration studied by Moser [26], Sitnikov [35] and Mathlouthi [21]. We consider two mass points $m_1 = m_2 > 0$, moving in the plane under Newton's attraction low in the elliptic orbits such that their center of mass O is at rest. We consider a third mass point m moving on the line perpendicular to the plane containing m_1 and m_2 , and going through O. We also suppose that mdoes not influence the motion of m_1 and m_2 . Let z be the coordinate describing the motion of m, so that z = 0 corresponds to O. The restricted three-body problem consists on determining z such that

$$\ddot{z} + V'(z) = 0.$$
 (1)

where

$$V(z) = -\frac{1}{(z^2 + r^2)^{\alpha/2}}.$$
 $0 < \alpha \le 1,$

and $r(t), t \in \mathbb{R}$ is the distance from the center of mass to anyone of the first two mass points. In the papers [21], [26] and [35], it was considered that

$$r(t) = \frac{1}{2}(1 - \varepsilon cost) + O(\varepsilon^2),$$

 $r : \mathbb{R} \longrightarrow \mathbb{R}$ is continuous, *T*-periodic, T > 0, and r(t) > 0, for all $t \in \mathbb{R}$, and the existence of a periodic solution was proved. In this work, we consider the autonomous case which corresponds to $\varepsilon = 0$, and we take r(t) = r > 0. So that m_1 and m_2 will have circular trajectories of constant radius r.

Definition 2.1. (Parabolic orbit)

A solution z of problem 1 is said to be a parabolic orbit (or homoclinic to infinity) if it is of class C^2 on \mathbb{R} and satisfies

(i)
$$\lim_{|t| \to +\infty} ||z(t)|| = \infty,$$

(ii)
$$\lim_{|t| \to +\infty} ||\dot{z}(t)|| = 0.$$
 (2)

The main result of this paper is the following

Theorem 2.1. For $0 < \alpha \leq 1$, problem 1 has at least one parabolic orbit.

3. Variational formulation

3.1. Step1. In this paragraph, we treat the case $0 < \alpha < 1$. For this, we consider the initial conditions

$$z(-mT) - z(mT) = \dot{z}(-mT) - \dot{z}(mT) = 0.$$

where T > 0 and r > 0. We denote by

$$H^1_{2mT} = \left\{ z \in H^1_{loc}(\mathbb{R}, \mathbb{R}) \quad ; z(-mT) = z(mT) \right\},$$

the space equipped with norm

$$||z||^{2} = \int_{-mT}^{mT} \left(|z(t)|^{2} + |\dot{z}(t)|^{2} \right) dt,$$

and we define on H_{2mT}^1 the functional

$$f_m(z) = \int_{-mT}^{mT} \left(\frac{1}{2} |\dot{z}(t)|^2 + \frac{1}{(z(t)^2 + r^2)^{\frac{\alpha}{2}}} \right) dt.$$

It's easy to prove that $f_m \in C^2(H^1_{2mT}, \mathbb{R})$ and to see that the solutions of the system are the critical points of the functional f_m . On the other hand, Palais-Smale condition $(PS)_c$ holds at c > 0. We recall the following result given in [21]

Lemma 3.1. a) The functional f_m is even and satisfies the Palais-Smale condition at every level c > 0.

b) If z is a solution of (1), then z(-t), -z(t) and -z(-t) are too.

Let Σ denote the class of sets $A \subset H^1_{2mT} \setminus \{0\}$ such that A is closed in H^1_{2mT} and symmetric with respect to 0. In what follows we define:

Definition 3.1. (Cogenus and Genus)

For $A \in \Sigma$ we define the cogenus of A to be

$$\gamma^{-}(A) = \inf \left\{ n \in \mathbb{N} / \exists \phi : A \to S^{n-1} \quad \text{odd and continuous} \right\},\$$

and the genus to be

 $\gamma^+(A) = \sup \left\{ n \in \mathbb{N} / \exists \phi : S^{n-1} \to A \text{ odd and continuous} \right\}.$

When such a ϕ doesn't exist, we set $\gamma^{\pm}(A) = \infty$

We recall in this proposition some of the properties of the cogenus. For more details, the reader can be referred, for example, to [4], [7] and [12].

Proposition 3.1. (i) If $z \neq 0$, $\gamma^{\pm}(\{z\} \cup \{-z\}) = 1$.

(ii) If Z is a k-dimensional subspace of H^1_{2mT} and $S = \{z \in Z/||z|| = r\}, r > 0,$ then, $\gamma^{\pm}(S) = k$.

(iii) If Z is a subspace of H^1_{2mT} of codimension k and $\gamma^{\pm}(A) > k$, then $A \cap Z \neq \emptyset$.

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For all $k \in \mathbb{N}$, we set

$$\gamma_k = \{A \in \Sigma, A \text{ is compact}, \gamma^-(A) \ge k\},\$$

and

$$c_k(m) = \inf_{A \in \gamma_{k+1}} \sup_{z \in A} f_m(z).$$

Then we have

 $c_k(m) \le c_{k+1}(m).$

Remark 3.1. (i) If f_m satisfies the Palais-Smale condition at the level $c_k(m)$, then, using the deformation lemma, we prove that $c_k(m)$ is a critical level of f_m .

(ii) If $c_k(m) = c_{k+1}(m) < f_m(0)$ and f_m satisfies the Palais-Smale condition at the level $c_k(m)$, then $f_m^{-1}(\{c_m\})$ contains infinitely many distinct critical points. (iii) $0 = c_0(m) < c_1(m)$.

(*iv*)
$$c_{2k_1}(m) < f_m(0)$$
, where $k_1 = \max\left\{k \in \mathbb{N}/k < \frac{mT}{\pi r}\right\}$.

For more details about these remarks, the reader can be referred, for example, to [4], [7], [12] and [21] and the references therein.

Next, we set

$$c_m = c_1(m) = \inf_{A \in \gamma_2} \sup_{z \in A} f_m(z).$$

We obtain the following result

Lemma 3.2. There exist two constants c > 0 and $0 < \theta < 1$, independent of m such that

$$c_m \le cm^{\theta}$$

Proof. Let $\frac{1}{2} < \beta < 1$, we consider the subspace of H_{2mT}^1 generated by the 2mT – periodic functions

$$\varphi_{\beta}(t) = |t|^{\beta}, \quad \forall t \in [-mT, mT]$$

 $\psi_{\beta}(t) = (mT - t)^{\beta}, \quad \forall t \in [0, mT], \text{ and even}$

and we set

$$K = \left\{ z \in H_{2mT}^1 \; ; \; z = \varepsilon_1 \tau \varphi_\beta + \varepsilon_2 (1 - \tau) \psi_\beta, \; \tau \in [0, 1] \text{ and } \varepsilon_1, \varepsilon_2 \in \{-1, 1\} \right\}$$

The set K is compact, symmetric not containing zero, homomorphic to the subset S^1 of H^1_{2mT} . This implies that the cogenus of K is equal to 2. Then,

$$c_m \le \max_{z \in K} f_m(z)$$

Indeed, we have

$$\dot{z}_{\tau,\varepsilon,\varepsilon_{1}} = \begin{cases} \beta \varepsilon \tau t^{\beta-1} - \beta (1-\tau)\varepsilon_{1}(mT-t)^{\beta-1} & \text{if} \quad t \in \left]0, mT\right[, \\ -\beta \varepsilon \tau t^{\beta-1} + \beta (1-\tau)\varepsilon_{1}(mT-t)^{\beta-1} & \text{if} \quad t \in \left]-mT, 0\right[\end{cases}$$

Then

$$\begin{split} \int_{-mT}^{mT} & |\dot{z}_{\tau,\varepsilon,\varepsilon_1}|^2 dt & \leq \quad (mT)^{2\beta-1} \left[\frac{2\beta^2}{2\beta-1} (\tau^2 + (1-\tau)^2) + 4\beta(1-\tau)\tau \frac{\Gamma(2\beta)}{\Gamma(\beta)^2} \right] \\ & \leq \quad (mT)^{2\beta-1} \left[\frac{2\beta^2}{2\beta-1} + \beta \frac{\Gamma(2\beta)}{\Gamma(\beta)^2} \right], \end{split}$$

where
$$\Gamma(t) = \int_0^{+\infty} s^{t-1} e^{-s} ds$$
. Moreover, we have
$$\int_0^{mT} \frac{dt}{\left(z_{\tau,\varepsilon,\varepsilon_1}^2(t) + r^2\right)^{\frac{\alpha}{2}}} \le (mT)^{1-\alpha\beta} \int_0^1 \frac{dt}{|\tau t^\beta - (1-\tau)(1-t)^\beta|^{\alpha}}.$$

For $\tau \in (0, 1)$, we have

$$\int_{0}^{\frac{1}{2}} \frac{1}{|\tau t^{\beta} - (1 - \tau)(1 - t)^{\beta}|^{\alpha}} dt \leq \\ \leq \begin{cases} \frac{2^{\alpha\beta}}{\beta} \left[\frac{(1 - \tau)^{-\alpha + 1} - (1 - 2\tau)^{-\alpha + 1}}{\tau (1 - \alpha)} \right], & \text{if } 0 < \tau \leq \frac{1}{2}. \\ \frac{2^{\alpha\beta}}{\beta} \left[\frac{(1 - \tau)^{-\alpha + 1} + (-1 + 2\tau)^{-\alpha + 1}}{\tau (1 - \alpha)} \right], & \text{if } \frac{1}{2} \leq \tau < 1. \end{cases}$$

In the other hand,

$$\int_{\frac{1}{2}}^{1} \frac{dt}{|\tau t^{\beta} - (1 - \tau)(1 - t)^{\beta}|^{\alpha}} \leq \\ \leq \begin{cases} \frac{2^{\alpha\beta}}{\beta} \left[\frac{\tau^{-\alpha+1} + (1 - 2\tau)^{-\alpha+1}}{(1 - \tau)(1 - \alpha)} \right], & \text{if } 0 < \tau \leq \frac{1}{2}. \\ \frac{2^{\alpha\beta}}{\beta} \left[\frac{\tau^{-\alpha+1} - (-1 + 2\tau)^{-\alpha+1}}{(1 - \tau)(1 - \alpha)} \right], & \text{if } \frac{1}{2} \leq \tau < 1. \end{cases}$$

Then, by Taylor expansion in 0 and 1, we deduce that

$$\int_0^{mT} \frac{dt}{\left(z_{\tau,\varepsilon,\varepsilon_1}^2(t) + r^2\right)^{\frac{\alpha}{2}}} \le c(mT)^{1-\alpha\beta},$$

where c is a constant depending on β and α . Then, taking

$$\theta = max(2\beta - 1, 1 - \alpha\beta),$$

we deduce that for all $\tau \in]0,1[$, we have $c_m \leq cm^{\theta}$. Hence the lemma follows. \Box

If $z_{\alpha,m}$ is a critical point corresponding to the critical level c_m , then we have Lemma 3.3. $||z_{\alpha,m}||_{\infty} \longrightarrow +\infty$ when $m \longrightarrow +\infty$.

Proof. We have,

$$\frac{2mT}{(\|z_{\alpha,m}\|_{\infty}^2 + r^2)^{\frac{\alpha}{2}}} = \int_{-mT}^{mT} \frac{dt}{(\|z_{\alpha,m}\|_{\infty}^2 + r^2)^{\frac{\alpha}{2}}} \le f_m(z_{\alpha,m}) \le cm^{\theta}.$$

This implies

$$\|z_{\alpha,m}\|_{\infty}^{2} \geq \frac{1}{c}m^{\frac{2(1-\theta)}{\alpha}} - r^{2} \longrightarrow +\infty \quad when \quad m \longrightarrow +\infty.$$

Remark 3.2. If $z_{\alpha,m}$ is a solution of (1), then $z_{\alpha,m}$ has at least one zero t_{0m} in [-mT, mT]. Then, there exists $n(m) \in \mathbb{Z}$ such that $n(m)T \leq t_{0m} < (n(m) + 1)T$. We set

$$\bar{z}_{\alpha,m}(t) = z_{\alpha,m}(t+n(m)T),$$

then $\bar{z}_{\alpha,m}$ is a solution of 1, having a zero in [0,T]. In the following, for simplicity, we denote $\bar{z}_{\alpha,m}$ by $z_{\alpha,m}$.

Lemma 3.4. $(\dot{z}_{\alpha,m})_{m\in\mathbb{N}}$ is uniformly bounded.

Proof. Let $m \in \mathbb{N}$ and $z_{\alpha,m}$ be a solution of 1. We take the energy function

$$E_m(t) = \frac{1}{2} |\dot{z}_{\alpha,m}(t)|^2 - \frac{1}{\left(z_{\alpha,m}^2 + r^2\right)^{\frac{\alpha}{2}}}.$$
(3)

The derivative of E_m with respect to t is such that $E'_m(t) = 0$ for all $t \in \mathbb{R}$ and for each solution $z_{\alpha,m}$ of (1). This implies that $E_m(t) = C(m)$.

We consider an extremum t_m of $z_{\alpha,m}$. Then we have

$$\dot{z}_{\alpha,m}(t_m) = 0.$$

So that we have

$$|C(m)| = \frac{1}{(z_{\alpha,m}^2(t_m) + r^2)^{\frac{\alpha}{2}}} \le \frac{1}{r^{\alpha}}$$

From 3 we deduce

$$|\dot{z}_{\alpha,m}(t)|^2 \le \frac{2}{\left(z_{\alpha,m}^2(t) + r^2\right)^{\frac{\alpha}{2}}} + 2|C(m)| \le \frac{4}{r^{\alpha}}$$

PROOF OF THE THEOREM FOR $0 < \alpha < 1$

As $z_{\alpha,m}$ vanishes on [0,T], for all $m \in \mathbb{N}$, Lemma 3 implies that $(z_{\alpha,m})_{m\in\mathbb{N}}$ is uniformly bounded on any compact set of \mathbb{R} . From 3.4, we deduce that $(z_{\alpha,m})_{m\in\mathbb{N}}$ is uniformly equicontinues on \mathbb{R} . Then applying the Ascoli-Arzelà theorem, we can extract a subspace of $(z_{\alpha,m})_{m\in\mathbb{N}}$ converging uniformly on any compact set of \mathbb{R} to a limit z_{α} , which is a solution of (1). z_{α} is also of class C^2 on \mathbb{R} . On the other hand, it satisfies the energy equation

$$E = \frac{1}{2} |\dot{z}_{\alpha}|^2 - \frac{1}{(z_{\alpha}^2 + r^2)^{\frac{\alpha}{2}}} = 0$$

Which implies that

$$\frac{1}{2}|\dot{z}_{\alpha}|^{2} = \frac{1}{(z_{\alpha}^{2} + r^{2})^{\frac{\alpha}{2}}} > 0.$$
(4)

So we can deduce that z_{α} is strictly increasing or strictly decreasing on \mathbb{R} . We suppose that z_{α} is bounded. So that there exists a constant B > 0 such that $|z_{\alpha}(t)| < B$ for all $t \in \mathbb{R}$. By 4, there exists $\gamma > 0$ such that $|\dot{z}_{\alpha}(t)| > \gamma$, for all $t \in \mathbb{R}$, which implies:

$$|z_{\alpha}(t) - z_{\alpha}(0)| = |\int_{0}^{t} \dot{z}_{\alpha}(t)dt| > \gamma|t| \to \infty \quad when \quad t \to \infty$$

This result contradicts the hypothesis " z_{α} is bounded". So that: $z_{\alpha}(t) \to \infty$ when $t \to \infty$.

And again from 4 we have: $|\dot{z}_{\alpha}(t)| \to 0$ when $|t| \to +\infty$. It follows that z_{α} is a parabolic orbit of problem 1.

3.2. Step2.

Lemma 3.5. For $0 < \alpha < 1$, the solution z_{α} vanishes on [0, T].

Proof. We suppose that z_{α} doesn't vanish on [0, T]. Than z_{α} doesn't change sign on [0, T]. Assume, for example, that $z_{\alpha}(t) > 0$, for all $t \in [0, T]$ (the case " $z_{\alpha}(t) < 0$, for all $t \in [0, T]$ " is similar). As z_{α} is continues on the compact set [0, T], there exists a constant $\beta > 0$ such that $z_{\alpha}(t) > \beta$ for all $t \in [0, T]$. But $z_{\alpha,m}$ converges uniformly to z_{α} on [0, T], so we deduce the existence of an integer N such that

$$m \ge N \Longrightarrow |z_{\alpha,m}(t) - z_{\alpha}(t)| < \frac{\beta}{2}$$

so that

$$z_{\alpha,m}(t) > \frac{\beta}{2} \quad \forall t \in [0,T] \quad \text{and} \quad \forall m \ge N$$

It follows that $z_{\alpha,m}(t) \neq 0 \quad \forall t \in [0,T]$. A contradiction.

PROOF OF THE THEOREM FOR $\alpha = 1$

Let $(\alpha_n)_{n \in \mathbb{N}}$ be an increasing sequence of]0,1[which converges to 1. For each $n \in \mathbb{N}$, we denote by $z_n = z_{\alpha_n}$ the solution, constructed below, of the problem

$$\begin{cases} \ddot{z_n}(t) + \alpha_n \frac{z_n(t)}{(z_n(t)^2 + r^2)^{\frac{\alpha_n}{2} + 1}} = 0, \\ \lim_{|t| \to \infty} |z_n(t)| = +\infty, \\ \lim_{|t| \to \infty} |\dot{z}_n(t)| = 0. \end{cases}$$
(5)

Applying 4 for z_n , we obtain

$$|\dot{z}_n(t)|^2 = \frac{2}{(z_n^2(t) + r^2)^{\frac{\alpha}{2}}} \le \frac{2}{r^{\alpha}} \quad \forall t \in \mathbb{R}.$$

which implies that $(\dot{z}_n)_n$ is uniformly bounded on \mathbb{R} .

As z_n vanishes on [0, T], for all $n \in \mathbb{N}$, we can deduce that $(\dot{z}_n)_n$ is uniformly bounded on every compact set of \mathbb{R} and uniformly equicontinues on \mathbb{R} . It follows that we can again apply the Ascoli-Arzelà Theorem (as we did with the sequence $(z_{\alpha,m})_{m\in\mathbb{N}}$ in Step 1) and extract from z_n a subsequence that is uniformly converging on any compact set of \mathbb{R} to a function z_1 which is a solution of problem 1 for $\alpha = 1$.

Arguing as in the previous step, we deduce by 4, as $|t| \rightarrow \infty$, that we have

$$\begin{cases} |z_1(t)| \longrightarrow \infty \\ \dot{z}_1(t) \longrightarrow 0 \end{cases}$$

So that z_1 is a parabolic orbit of problem 1.

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References

- V.M. ALEKSEEV, Quasirandom Dynamical Systems, I, II, III. Math. St., 5, 73-128 (1968); 6, 503-560 (1968); 7, 1-43 (1969).
- [2] A. AMBROSETTI, Critical Points and Nonlinear Variational Problems. Bull. Soc. Math. France, 120, (1992), n.49, Cours de la chaire Lagrange.
- [3] A. AMBROSETTI, V. COTI-ZELATI, Non Collision Periodic Solution for a Class of Symmetric 3-Body Type Problems, *Topological Method Nonlinear Anal*, 3, 197-207 (1994).
- [4] A. AMBROSETTI, P.H. RABINOWITZ, Dual Variational Methods in Critical Point Theory and Applications, J.Funct. Analysis, 14, 349-360 (1973).
- [5] A. BAHRI, P.H. RABINOWITZ, Periodic Solutions of Hamiltonian Systems of 3-Body Type, Ann. IHP, Anal. Non Lineaire, 8, 561-649 (1991).
- [6] A. BAHRI, B. D'ONOFRIO, Exponential Growth of the Number of Periodic Orbits for 3-Body Type Problems, *Maghreb Math. Rev.*, 1, n.1, 1-14 (1992).
- [7] A. BAHRI, P.L. LIONS, Morse Index of Min-Max Critical Points I. Applications to Multiplicity Results, comm. pure and application, 8, Vol XLI, 1027-1037 (1988).
- [8] F. BENCI, F. GIANNONI, A New Proof of the Existence of a Break Orbit, Advanced Topics in the Theory of Dynamical Systems (Trento 1987). Notes Rep. Math. Sci. Energ., Academic Press, 6, (1990).
- [9] V. BENCI, F. GIANNONI, Homoclinic Orbits on Compact Manifolds, Academic Press, Inc. January, 10, (1990).
- [10] U. BESSI, Multiple Homoclinic Orbits for Autonomous, Singular Potentials, SNS, (1990).
- [11] S.V. BOLOTIN, The Existence of Homoclinic Motions, Vestnik Moskow Univ. Ser 1 Math. Mekh., 6, 98-103 (1980).
- [12] L. CAKLOVIĆ, J.L. SHU, M. WILLEM, A Note on Palais-Smale Condition and Convexity, Differential and integral equations, 3, 799-800 (1990).
- [13] C.D. CLARK, A Variant of the Ljusternik-Schnirelmann Theory, Indiana Uni Math. J, 22, 391-419 (1969).
- [14] V. COTI-ZELATI, A Class of Periodic Solutions of the N-Body Problem, Cel. Mech. and Dyn. Astr., 46, 177-186 (1989).
- [15] V. COTI-ZELATI, Periodic Solutions of the N-Body Type Problem, Annales Inst. H. Poincaré Anal. Nonlin. 7, 477-192 (1990).
- [16] V. COTI-ZELATI, I. EKELAND, E. SÉRÉ, Solutions Doublement Asymptotiques de Systèmes Hamiltoniens Convexes. C.R.A.S, t. 310, Série I, 631-633 (1990).
- [17] V. COTI-ZELATI, I. EKELAND, E. SÉRÉ, A Variational Approach to Homoclinic Orbits in Hamiltonian Systems. Math. Annalen 288, 133-160 (1990).
- [18] G. DELL'ANTONIO, Classical Solutions for a Perturbed N-Body System, A. Matzeu (Ed), Topol. Nonlin. Anal. II, Birkhäuser Vignolied, 1-86 (1997).
- [19] I. EKELAND, Convexity Methods in Hamiltonian Mechanics, Springer Verlag (1990).
- [20] R. MC. GEHEE, Parabolic Orbits of the Restricted Three-Body Problem, Academic Press, New York and London, ASubsidiary of Hareourt Brace Joranovich, Publicher (1973).
- [21] S. MATHLOUTHI, Periodic Orbits of the Restricted 3-Body Problem, Trans. Am. Math.Sic., 350, n.6, 2265-2276, (1998).
- [22] P. MAJER, S. TERRACINI, Multiple Periodic Solutions to Some N-Body Type Problems Via Collision Index, preprint (1993).
- [23] P. MAJER, S. TERRACINI S. Periodic Solutions to Some Problems of N-Body Type, Arch. Rational. Mech. Anal. 1124, 381-404 (1993).
- [24] P. MAJER, S. TERRACINI, Periodic Solutions to Some Problems of N-Body Type, the Fixed Energy Case, Duke Math. J., 69, n.3, 683-697 (1993).
- [25] J. MAWHIN, M. WILLEM, Critical Point Theory and Hamiltonian Systems, Applied Mathematical Sciences, Springer Verlag, Berlin Heidelberg New York, 78 (1989).
- [26] J. MOSER, Stable and Random Motion in Dynamical Systems, Ann. Math. Studies. Princeton University Press, 77 (1973).
- [27] H. POINCARÉ, Les Méthodes Nouvelles de la Mécanique Céleste, Gauthier-Villars. Paris,(1899).
- [28] P.H. RABINOWITZ, Minimax Methods in Critical Point Theory with Applications to Differential Equations, C.B.M.S. Reg. Conf. Ser. in Math. Amer. Math. Soc. Providence, R.I., 65 (1986).

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- [29] H. RIAHI, Periodic Orbits for N-Body Type Problems: The Fixed Energy Case, Trans. Amer. Math. Soc., 347, n.1, 4663-4685 (1995).
- [30] H. RIAHI, Study of the Generalised Solutions of n-Body Type Problems with Week Force, Nonlinear Anal. Theory, Methods Appl. 28, n.1, 49-59 (1997).
- [31] H. RIAHI, Study of the Critical Points at Infinity Arising From the Failure of the Palais-Smale Condition for N-Body Type Problems, Mem. Amer. Math. Soc. (1999).
- [32] E. SÉRÉ, Existence of Infinitely Many Homoclinic Orbits in Hamiltonian Systems, Math. Zeitschrift, 209, 27-42 (1992).
- [33] E. SERRA, S. TERRACINI, Noncollision Periodic Solutions to some 3-Body like Problems, Preprint SISSA, Trieste (1990).
- [34] E. SERRA, S. TERRACINI, Collisionless Periodic Solutions to Some 3-Body Problems, Arch. Rational Mech. Anal., 120, 305-325 (1992).
- [35] K. SITNIKOV, Existence of Oscillating for the Three-Body Problem, J. Dokl. Acad Nauk USSR, 133, n.2, 303-306 (1960).
- [36] D.L. SLOTNICK, Asymptotic Behavior of solutions of Canonical Systems Near a Closed, Unstable Orbit, Ann. Math. Studies (English), 41, 85-110 (1958).
- [37] M. STRUWE, Variational Methods, Springer Verlag, (1990).
- [38] E. VITILLARO, Non Collision Periodic Solution of Fixed Energy for a Symmetric N-Body Type Problems, Variational and Local Methods in the Study of Hamiltonian Systems (Trieste, 1994), World Sci. Publishing, River Edge, NJ, 202-211 (1995).

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